




High-performance Ge/Si electro-absorption optical modulator up to 85 °C and its highly efficient photodetector operation

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Abstract: We studied a high-speed Ge/Si electro-absorption optical modulator (EAM) evanescently coupled with a Si waveguide of a lateral p–n junction for a high-bandwidth optical interconnect over a wide range of temperatures from 25 °C to 85 °C. We demonstrated 56 Gbps high-speed operation at temperatures up to 85 °C. From the photoluminescence spectra, we confirmed that the bandgap energy dependence on temperature is relatively small, which is consistent with the shift in the operation wavelengths with increasing temperature for a Ge/Si EAM. We also demonstrated that the same device operates as a high-speed and high-efficiency Ge photodetector with the Franz-Keldysh (F-K) and avalanche-multiplication effects. These results demonstrate that the Ge/Si stacked structure is promising for both high-performance optical modulators and photodetectors integrated on Si platforms.

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1. Introduction

Silicon photonics has recently attracted considerable attention because of its low cost, low power consumption, and high bandwidth for optoelectronic solutions for applications ranging from telecommunications to chip-to-chip interconnects [1–3]. To realize an effective photonic–electronics convergence system, it is very important to integrate a high-speed optical modulator with a Si-based optical circuit over a wide range of temperatures.

Among the various Si optical modulators (Si-MODs) demonstrated so far, Mach–Zehnder Si-MODs based on the free-carrier plasma dispersion effect have been mostly reported for high-speed modulation and broad-wavelength operation [4–11]. However, carrier-depletion Si modulators require a relatively long phase shifter or a high driving voltage because of the weak plasma dispersion effect in Si [12], and they are not favorable for large-scale integration. It has been reported that ring-resonator-type Si-MODs can realize high-speed and low-power operation. However, the operation wavelength must be controlled to an accuracy of less than 1 nm, which makes practical use difficult [13,14].

To achieve a low-power and high-density interconnect system, an optical modulator with a very small capacitance is required. The GeSi electro-absorption optical modulator (EAM) is promising because its electrical capacitance is approximately 10 fF and the device length is dozens of micrometers [15–19]. It has been reported that a Ge layer on a Si substrate has a tensile strain as large as 0.2%, which reduces the direct band-gap energy from 0.80 eV to 0.77 eV [20]. This leads to an L-band operation of EAMs using such a strained Ge layer. Recently, we have

reported that decreasing the Ge width enabled the EAM to operate in the shorter wavelength range of the C-band [21], although the shift in the operation wavelengths with temperature should be concerned because of the change in the direct bandgap energy. The operating wavelengths of a Ge/Si EAM have been reported to shift to longer wavelengths at high temperatures because the direct bandgap energy decreases with temperature [22]. In addition, a demonstration of a GeSi EAM and Ge photodetector (PD) integrated on a Si platform has been reported, which is promising for a high-density and low-power photonics-electronics convergence system [15–24].

In this paper, we study a high-speed Ge/Si EAM evanescently coupled with a Si waveguide of a lateral p–n junction for a high-bandwidth optical interconnect over a wide range of temperatures from 25 °C to 85 °C. We demonstrated C-band wavelength operation and 56 Gbps high-speed operation at high temperatures up to 85 °C. We also demonstrated a high-responsivity and wide-bandwidth Ge PD in the C + L band wavelengths with the same structure as the EAM. The Franz-Keldysh (F-K) and avalanche-multiplication effects realize the high responsivity with a lower V_{dc} compared with a report on a Ge PD of a vertical p–n junction [25].

2. Experiment and results

2.1. Ge/Si EAM

Figure 1(a) shows a schematic cross-section of a Ge/Si EAM on a Si rib waveguide with a lateral pn junction and cross-sectional TEM (transmission electron microscope) images of a Ge/Si EAM for Fig. 1(b) 0.3 μm -width, Fig. 1(c) 0.6 μm -width, and Fig. 1(d) 1.0 μm -width. The Ge/Si EAM consists of an evanescently coupled structure with a Si rib waveguide [21]. The stack Ge/Si layer width was changed from 0.3 μm to 1.0 μm , and the Ge height was about 300 nm.

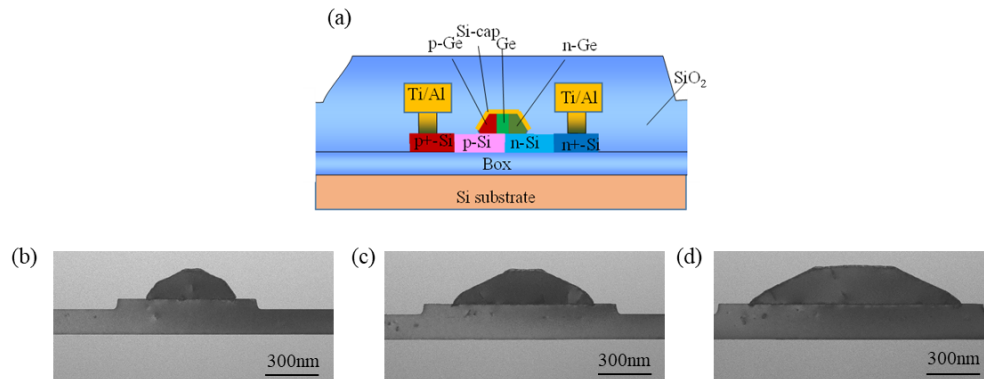


Fig. 1. (a) Schematic diagram of Ge/Si EAM. Cross-sectional TEM image of Ge/Si EAM for (b) 0.3 μm -width, (c) 0.6 μm -width, and (d) 1.0 μm -width.

Figure 2 schematically shows the Ge/Si EAM fabrication process. The fabrication process started from a 300 mm-diameter silicon-on-insulator (SOI) wafer with a 200-nm SOI thickness. Boron (B) and phosphorus (P) ions were implanted into the SOI layer, and the wafers were annealed to form a lateral p–n junction in the SOI layer. Then, a Si pedestal was patterned using immersion ArF lithography and dry etching. Subsequently, nominally 500 nm-thick epitaxial Ge was selectively grown on the Si pedestal using the ultra-high-vacuum chemical vapor deposition method. In the present case, the Ge thickness decreased from 500 nm to approximately 300 nm with a decrease in the Ge width to less than 1 μm because the Ge growth on the (001) crystalline plane is interfered with those on the (111) and (311) sidewall planes. In the Ge epitaxial growth process, we applied a 20–30 nm-thick $\text{Si}_{0.5}\text{Ge}_{0.5}$ buffer layer, followed by a pure Ge layer growth. After Ge epitaxial growth, a post-annealing process at approximately 800 °C was applied to

the wafer for 30 min to improve the crystallinity of the Ge layer on the SOI substrate. Then a 20 nm-thick Si-capping layer was deposited onto the Ge layer to passivate the Ge surface. Next, we implanted B and P ions at the Ge sidewalls with the doping density of $2 \times 10^{18}/\text{cm}^3$ and annealed the wafers to form a lateral p-i-n junction in the Ge layer. Subsequently, a SiO_2 upper-clad layer was deposited, and contact holes were formed by UV lithography and a dry-etching process. Finally, the metal electrodes of the Ti/TiN/Al layers were deposited and patterned.

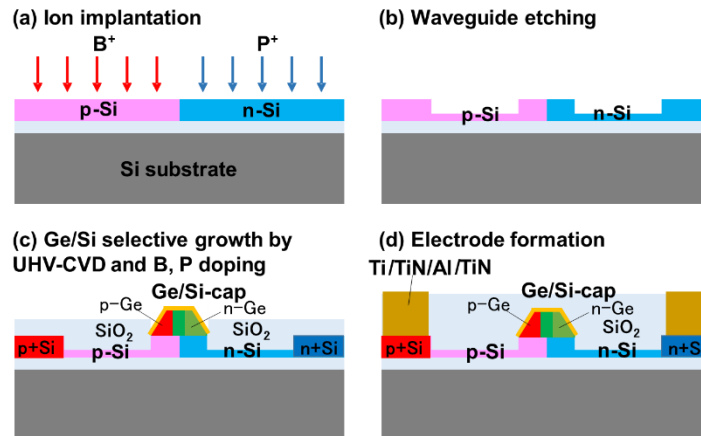


Fig. 2. Fabrication process flow of Ge/Si EAM and Ge/Si PD [21].

Figure 3 shows the frequency EO response for the Ge/Si EAM with 0.3- μm width and 40- μm length at room temperature [21]. Its bandwidth was more than 67 GHz at 2.0 V_{dc} , which is enough for high-speed operation. Figure 4 shows a schematic of the measurement system for the high-speed Ge/Si EAM. A continuous wave (CW) laser light source with a wavelength of 1550 nm and optical power of 13 dBm was applied to the measurement system. Polarization-controlled light was input into the Ge/Si EAM and amplified by an erbium-doped fiber amplifier (EDFA) to detect the optical signal using a sampling oscilloscope or a real-time oscilloscope. To drive the Ge/Si EAM at high speed, an arbitrary waveform generator (AWG) was used to generate high-speed electrical signals of 112 Gbps pulse-amplitude-modulation 4 (PAM-4) and 100 Gbps (non-return-to-zero) NRZ. The high-speed electrical signal to drive the Ge/Si EAM had 1.6 V_{pp} amplitude for the 112 Gbps PAM-4 and 0.8 V_{pp} amplitude for the 100 Gbps NRZ and a V_{dc} of 2.0 V was applied through the bias T. In case of measurement for the high-speed characteristics for the Ge/Si EAM at high temperature, we used the pulse pattern generator (PPG) with 2.5 V_{pp} amplitude for 56 Gbps NRZ and 2.0 V_{dc} .

Figure 5(a) and (b) show the output waveforms of the 112 Gbps PAM-4 and 100 Gbps NRZ with a 2^7-1 pseudo random binary sequence (PRBS) at room temperature. The extinction ratios were 3.1 dB for the 112 Gbps PAM-4 with a 1.6 V_{pp} RF applied voltage and 1.5 dB for the 100 Gbps NRZ with 0.8 V_{pp} RF applied voltage. We demonstrated clear eye opening both for the 112 Gbps PAM-4 and 100 Gbps NRZ modulation formats, which shows that the Ge EAM is promising for high-speed optical interconnects.

Figure 6 shows the experimental results of the optical transmission dependence on the applied voltage at Fig. 6(a) 25 $^{\circ}\text{C}$, Fig. 6(b) 45 $^{\circ}\text{C}$, Fig. 6(c) 65 $^{\circ}\text{C}$, and Fig. 6(d) 85 $^{\circ}\text{C}$ for 0.3- μm -wide and 40- μm -long Ge/Si EAMs. The optical transmission for the Ge/Si EAM is defined as the difference between the optical transmission power of a reference Si channel waveguide and a Si channel waveguide with a Ge/Si EAM in the same chip. It includes the optical coupling loss between a Si WG and Ge/Si EAM. When increasing the reverse bias voltage for a 0.3- μm wide and 40- μm -long Ge/Si EAM at 25 $^{\circ}\text{C}$, the optical transmission power decreased from -6

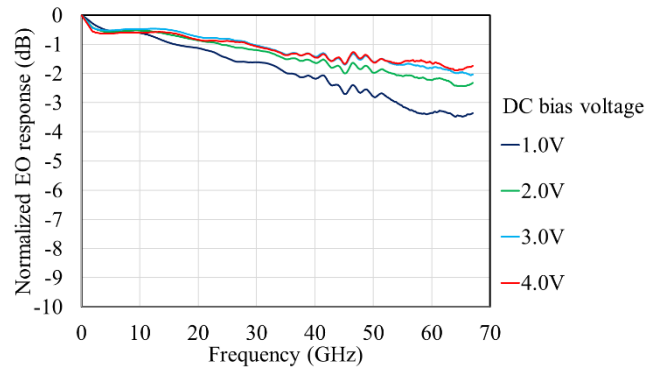


Fig. 3. Frequency EO response for Ge/Si EAM with 0.3- μm width and 40- μm length [21].

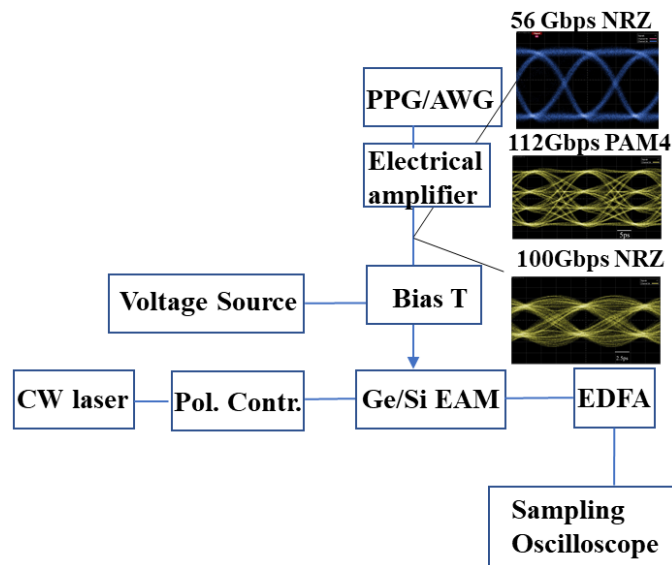


Fig. 4. Schematic diagram of high-speed measurement system for Ge/Si EAM with 0.3- μm width and 40- μm length.

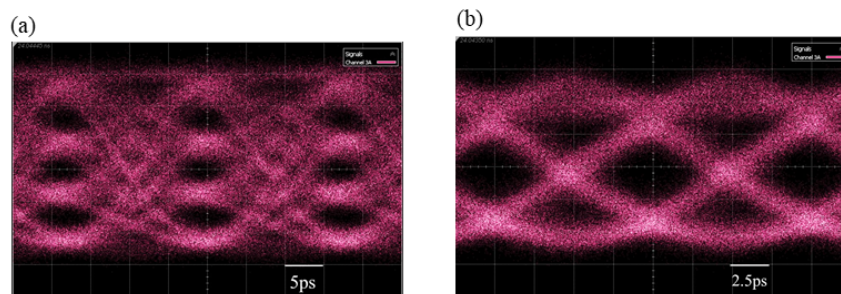


Fig. 5. Output waveforms of (a) 112 Gbps PAM-4 with 2^7-1 PRBS and (b) 100 Gbps NRZ with 2^7-1 PRBS at 1550 nm wavelength, measured at room temperature.

to -9 dB at around a 1550-nm wavelength, which originates from the F-K effect. On the other hand, a 1.0 μm -wide Ge/Si EAM operated at wavelengths between 1580 and 1600 nm, which is consistent with the tensile-strained Ge bandgap energy on Si [19]. The optical transmission spectra for 0.3 μm -wide and 0.6 μm -wide Ge/Si EAMs shifted to shorter wavelengths than that of a 1.0- μm -wide GeSi EAM. This behavior is because the tensile-strain in the smaller-width of a Ge layer is elastically relaxed [26]. From the I-V characteristics, the leakage current of the GeSi EAM is less than 1 μA , which does not affect the optical transmission loss.

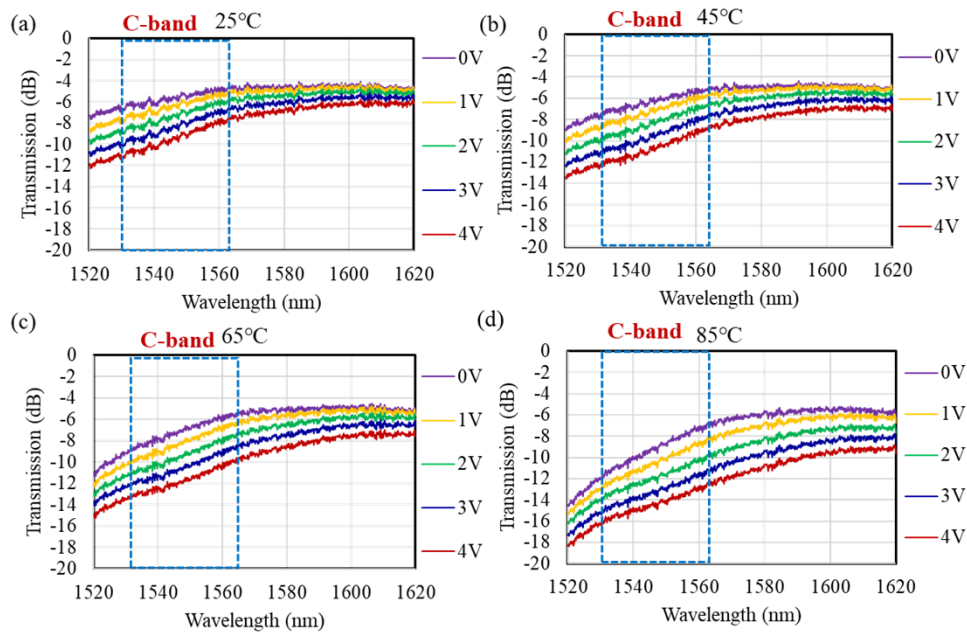


Fig. 6. Experimental results of optical transmission spectrum dependence on V_{dc} for (a) 25°C, (b) 45°C, (c) 65°C, and (d) 85°C.

Figures 7(a) and (b) shows optical transmission spectra at different temperatures and the wavelength at the transmission of -20 dB transmission as a function of temperature, respectively. In this study, absorption edge wavelength was defined as that of optical transmission at -20 dB. With increasing the temperature, the optical absorption edge wavelength shifts to longer wavelength, because direct bandgap energy decreases with temperature [22]. The slope of absorption edge wavelength dependence on temperature was 0.60 nm/°C, which is smaller than that for the previously reported GeSi modulator [16]. Theoretical analysis shows that an increase of Si composition in a GeSi alloy decreases the direct band gap energy and its temperature dependence a little [22]. Therefore, GeSi alloying would affect the temperature dependence of the optical absorption-edge wavelength.

Figure 8 shows photoluminescence (PL) spectrum-peak wavelength dependence on temperature for a bulk Ge layer, a strained Ge layer on Si, and a 0.6 μm -width Ge/Si EAM. In this measurement, an excitation laser wavelength was 785 nm. As for a bulk Ge layer, the slope of PL spectrum peak wavelength vs. temperature was 0.76 nm/°C, which is comparable with the theoretical value [22]. On the other hand, a 0.6 μm -width of Ge/Si EAM shows smaller dependence of PL spectrum peak wavelength on temperature. The slope of the PL spectrum-peak wavelength dependence is 0.54 nm/°C, which is smaller than the theoretical value for $\text{Ge}_{0.915}\text{Si}_{0.085}$. Therefore, GeSi alloying with various GeSi compositions would contribute to smaller temperature dependence

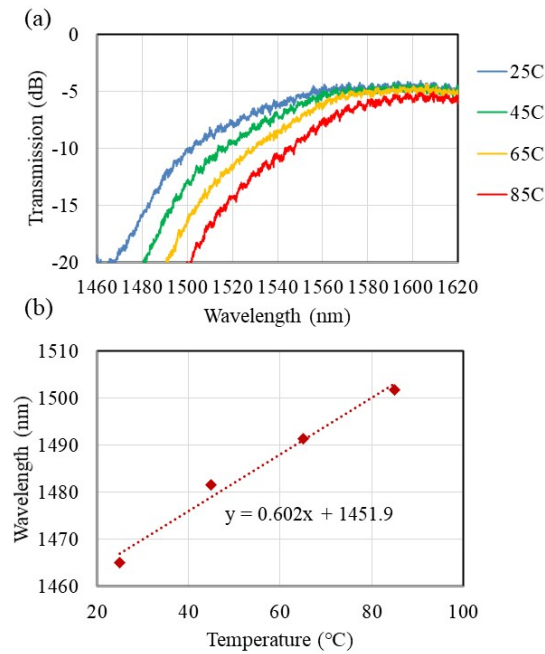


Fig. 7. (a) Experimental result of optical transmission spectrum dependence on temperature and (b) measured wavelength at -20 dB transmission for different temperatures for a GeSi EAM.

of bandgap energy, that is, operation wavelength dependence on temperature for a Ge/Si EAM would decrease with GeSi alloying at the Ge/Si interface.

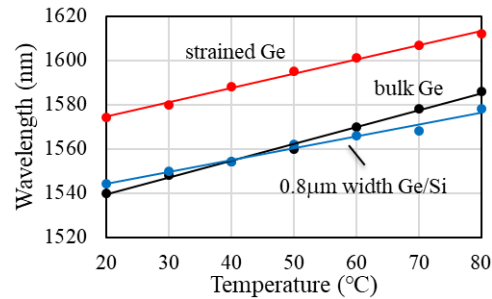


Fig. 8. Photoluminescence-spectrum-peak wavelength dependence on temperature for a bulk Ge layer, a strained Ge layer on Si, and a 0.6 μm-wide Ge/Si EAM.

We analyzed the Ge/Si EAM using Raman spectroscopy with an excitation laser wavelength of 457 nm to investigate the crystalline strain and GeSi alloying of the Ge/Si EAM. From the Raman spectrum, we can verify the GeSi mixed crystalline layer in the Ge/Si stacked structure, in addition to the Ge layer [27,28]. The Raman spectrum peak of Ge-Ge bonding was approximately 298 cm^{-1} , which is consistent with that of a tensile-strained Ge layer. In addition, a small broad peak originating from the GeSi mixed-crystalline structure was observed. Therefore, a GeSi mixed crystalline layer was formed by the Ge epitaxial growth process, which would contribute to the smaller temperature dependence of the operating wavelengths [21,22].

Figure 9(a)–(d) show the output waveforms at 56 Gbps with $2^{31}-1$ PRBS (pseudo-random bit sequence) for a $1.55\ \mu\text{m}$ wavelength at $2.0\ \text{V}_{\text{dc}}$ for a $40\ \mu\text{m}$ -long Ge/Si EAM at $25\ ^\circ\text{C}$, $45\ ^\circ\text{C}$, $65\ ^\circ\text{C}$, and $85\ ^\circ\text{C}$. The insertion loss was $5.5\ \text{dB}$, and ER was $3.0\ \text{dB}$ at room temperature for the Ge/Si device at $25\ ^\circ\text{C}$. In this experiment, optical input power was $3\ \text{dBm}$ via a lensed optical fiber and optical coupling loss between a lensed optical fiber and the Si waveguide spot size converter was about $-2.5\ \text{dB}$. RF drive voltage was $2.5\ \text{V}_{\text{pp}}$ and V_{dc} was $2.0\ \text{V}$ at each temperature. To observe the high-speed output waveform, EDFA was used to maintain the time-averaged output optical power to be $-1.2\ \text{dBm}$ for each temperature. A clear eye opening was obtained at temperatures up to $85\ ^\circ\text{C}$. Extinction ratios were $3.0\ \text{dB}$ at $25\ ^\circ\text{C}$, $3.2\ \text{dB}$ at $45\ ^\circ\text{C}$, $3.0\ \text{dB}$ at $60\ ^\circ\text{C}$, and about $2.0\ \text{dB}$ at $85\ ^\circ\text{C}$. The optical losses at $25\ ^\circ\text{C}$ and $85\ ^\circ\text{C}$ at $1550\ \text{nm}$ wavelength were $-7\ \text{dB}$ and $-12\ \text{dB}$ at $2.0\ \text{V}_{\text{dc}}$. At $85\ ^\circ\text{C}$, the optical coupling between a lensed fiber and a Si waveguide for a Ge/Si EAM is unstable in the measurement system, which increases the optical intensity jitter in the output waveform. The frequency bandwidth was greater than $67\ \text{GHz}$ for $40\ \mu\text{m}$ length of a Ge/Si EAM. The electrical capacitance was estimated to be approximately $10\ \text{fF}$ for the $40\ \mu\text{m}$ length. Therefore, the Ge/Si EAM is promising for low-power and high-bandwidth optical interconnects by improving the insertion optical loss.

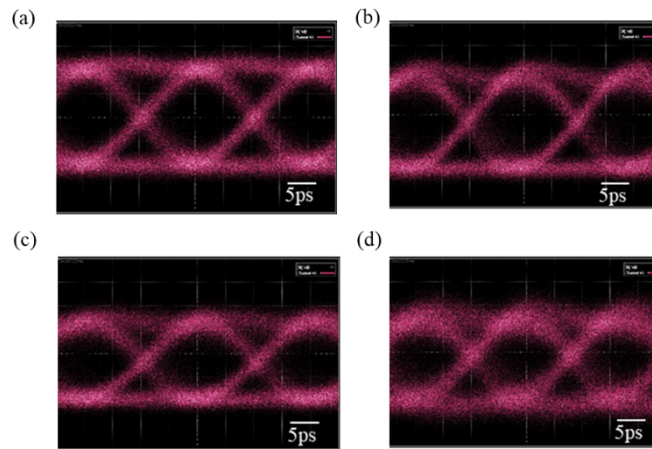


Fig. 9. Output waveforms at 56 Gbps with $2^{31}-1$ PRBS for $1.55\ \mu\text{m}$ wavelength at (a) $25\ ^\circ\text{C}$, (b) $45\ ^\circ\text{C}$, (c) $65\ ^\circ\text{C}$, and (d) $85\ ^\circ\text{C}$ for Ge/Si EAM.

3. Ge/Si photodetector

3.1. High-photoresponsivity by F-K effect

Next, we studied the photoresponsivity of the Ge–Si stacked structure. Figure 10(a) and (b) show the I–V characteristics and OE response dependence on the frequency and V_{dc} for the $0.3\ \mu\text{m}$ wide and $40\ \mu\text{m}$ -long Ge/Si stacked structure. In this experiment, optical input power was $3\ \text{dBm}$ via a lensed optical fiber and optical coupling loss between a lensed optical fiber and the Si waveguide spot size converter was about $-2.5\ \text{dB}$. With an increase in V_{dc} , the photoresponsivity increased up to approximately $0.83\ \text{A/W}$ with $4.0\ \text{V}_{\text{dc}}$ at $1550\ \text{nm}$ wavelength. The $3\ \text{dB}$ bandwidth in the OE frequency response was $60\ \text{GHz}$ for more than $2\ \text{V}_{\text{dc}}$.

Figure 11(a), (b), and (c) show the photoresponsivity dependences on the input light wavelength with increasing V_{dc} for the $0.3\ \mu\text{m}$, $0.6\ \mu\text{m}$, and $1.0\ \mu\text{m}$ -wide Ge/Si stacked structures of $40\ \mu\text{m}$ -length. The photoresponsivity enhancement was obtained by the F-K effect with an increase in V_{dc} . As for the $0.3\ \mu\text{m}$ -wide Ge/Si stacked structure, photoresponsivity enhancement by the F-K effect was observed for a large wavelength bandwidth from $1500\ \text{nm}$ to $1620\ \text{nm}$. In this

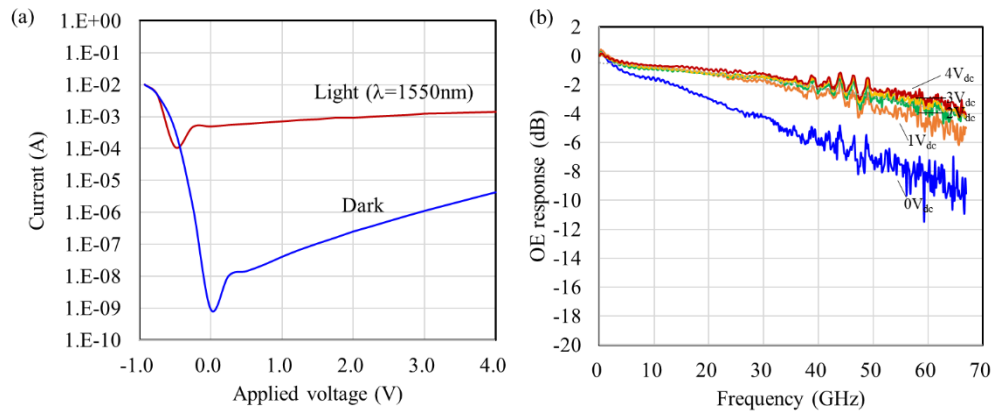


Fig. 10. (a) I-V characteristics with and without $1.55\ \mu\text{m}$ light. (b) Frequency dependence of OE response with 0 to $4\ V_{dc}$.

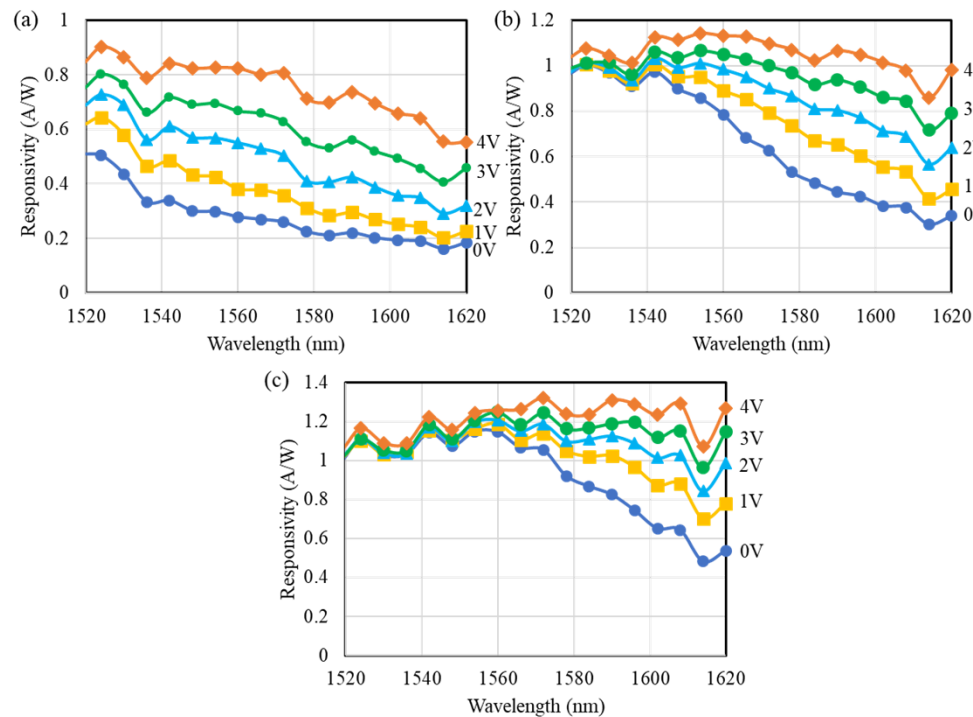


Fig. 11. Bias voltage (V_{dc}) dependence of photoresponsivity for (a) $0.3\ \mu\text{m}$ -width, (b) $0.6\ \mu\text{m}$ -width, and (c) $1.0\ \mu\text{m}$ -width Ge/Si structure of $40\ \mu\text{m}$ -length.

experiment, optical input power was 3dBm via a lensed optical fiber and optical coupling loss between a lensed optical fiber and the Si waveguide spot size converter was measured from 1500 nm to 1620 nm wavelengths. To calculate photoresponsivity, input optical power to the Ge/Si photodetector was calibrated based on the optical coupling loss dependence on wavelength. The wavelength bandwidths with photoresponsivity enhancement by F-K effect shift to longer wavelength for the $0.6\ \mu\text{m}$ and $1.0\ \mu\text{m}$ wide Ge/Si structures. This is because the bandgap in Ge decreases owing to the tensile strain caused by the difference in the thermal expansion coefficients

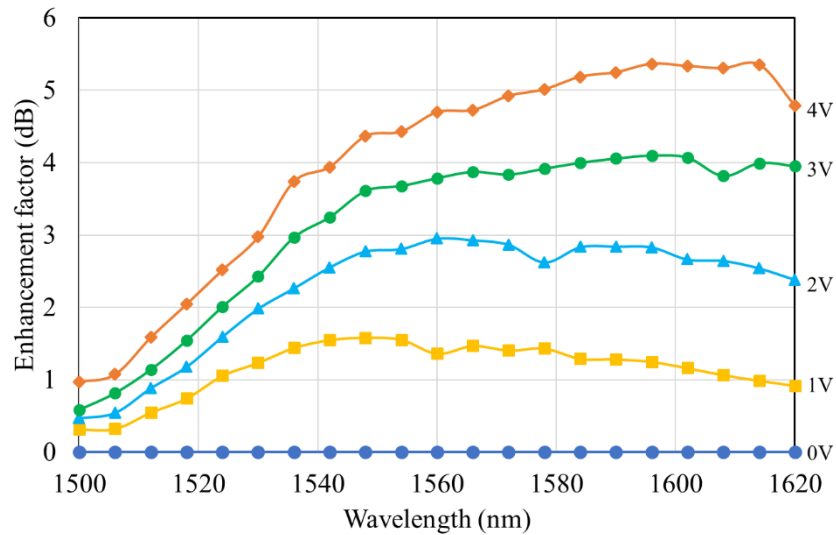


Fig. 12. Photoresponsivity enhancement factor dependence on V_{dc} for 0.3 μm -wide Ge/Si stacked structure of 40 μm -length.

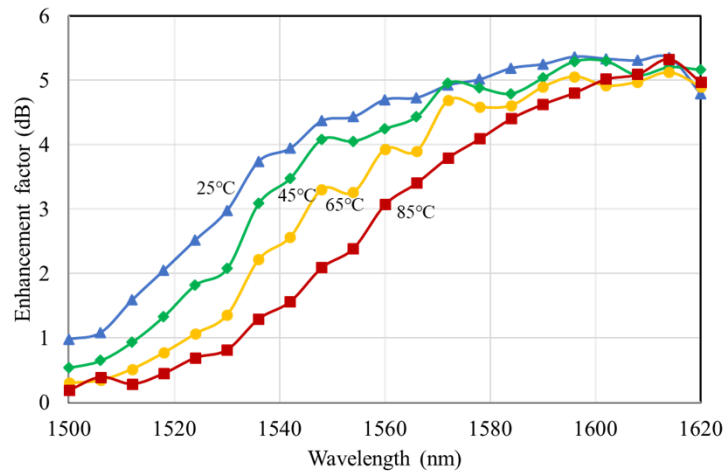


Fig. 13. Photoresponsivity enhancement factor dependence on temperature at $4V_{dc}$ for 0.3 μm -wide Ge/Si stacked structure of 40 μm -length.

of the Ge and Si layers [26]. The smaller width of the Ge/Si stacked structures shows relaxation in tensile strain because of the decrease in the contact area at the Ge/Si stacked interface [26].

Figure 12 shows the enhancement factor of the photoresponsivity dependence on wavelength with an increase in V_{dc} for the 0.3 μm -width Ge/Si stack structure. About 5 dB enhancement of photocurrent could be obtained for larger wavelengths more than 1560 nm. Therefore, the Ge/Si stacked structure with 0.3 μm -width shows high-photoresponsivity, high speed, and large bandwidth in the operating wavelengths.

Figure 13 shows the enhancement factor of the photoresponsivity dependence on wavelength with an increase in temperature at 4 V_{dc} for the 0.3 μm -width Ge/Si stack structure. With increase in temperature, optical absorption edge wavelength shift to longer wavelength [29], and photoresponsivity enhancement factor by F-K effect was also observed in the longer wavelength.

Optical absorption edge wavelength dependence on temperature was similar with that shown in Fig. 7(b).

3.2. High-photoresponsivity by the avalanche-multiplication effect

The avalanche-multiplication effect was studied to obtain a higher photoresponsivity in addition to that of the F-K effect [23]. Figure 14(a) and (b) show the I-V characteristics with and without 1.55 μm light and the photoresponsivity dependence on wavelengths with an increase in V_{dc} for a 0.3 μm -wide and 40 μm -long Ge/Si stacked structure. From the I-V characteristics, current amplification by avalanche phenomena in the Ge layer was observed at approximately 5 V_{dc} .

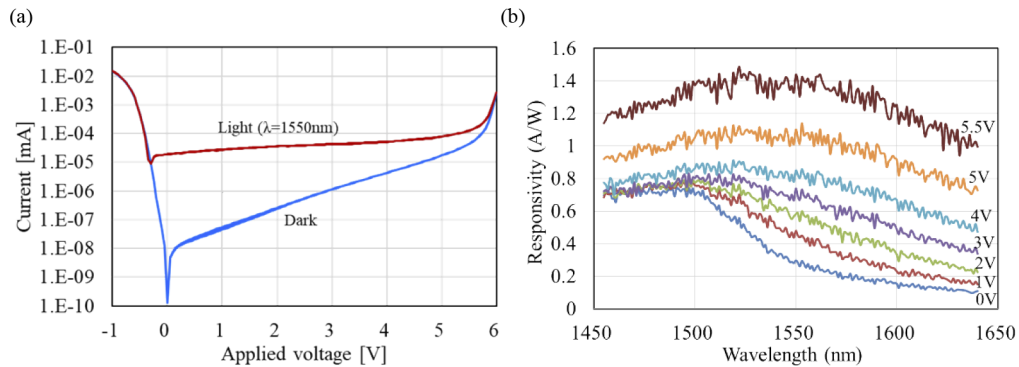


Fig. 14. (a) I-V characteristics with and without 1.55 μm light. (b) Photoresponsivity dependence on wavelengths with increase in V_{dc} for 0.3 μm -wide and 40 μm -long Ge/Si stacked structure.

As shown in Fig. 13(b), the photoresponsivity enhancement by the avalanche effect was measured over a large range of wavelengths greater than 5 V_{dc} . This is different from that of the F-K effect, because the F-K effect is observed mainly around the direct bandgap energy, which is higher than 1500 nm for this device. The multiplication factor of the avalanche multiplication-effect is estimated to be approximately 2. By using the F-K effect and avalanche multiplication-effect, a 7 dB enhancement of photoresponsivity was obtained at a wavelength of 1550 nm.

4. Conclusion

We studied a high-speed Ge/Si EAM evanescently coupled with a Si waveguide of a p-n junction for a high-bandwidth optical interconnect over a wide range of temperatures from 25 $^{\circ}\text{C}$ to 85 $^{\circ}\text{C}$. We demonstrated 56 Gbps high-speed operation at high temperatures up to 85 $^{\circ}\text{C}$. From the photoluminescence spectra, we confirmed that the bandgap energy dependence on temperature is relatively small, which is consistent with that of the operation wavelengths with increasing temperature for a Ge/Si EAM. We also demonstrate a high-speed and high-efficiency Ge photodetector. The photoresponsivity was enhanced by the F-K effect and avalanche-multiplication effect in the C + L band wavelengths. These results demonstrate that the Ge/Si stacked structure is promising for both high-performance optical modulators and photodetectors integrated on Si platforms.

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Data availability. Data underlying the results presented in this paper are not publicly available at this time but may be obtained from the authors upon reasonable request.

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